

EXISTENCE OF STATIONARY TURBULENT FLOWS WITH VARIABLE POSITIVE VORTEX INTENSITY

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ABSTRACT. We prove the existence of stationary turbulent flows with arbitrary positive vortex circulation on non simply connected domains. Our construction yields solutions for all real values of the inverse temperature with the exception of a quantized set, for which blow-up phenomena may occur. Our results complete the analysis initiated in [21].

1. INTRODUCTION AND MAIN RESULTS

Motivated by the statistical mechanics description of turbulent 2D Euler flows in equilibrium, we are interested in the existence of solutions to the following problem:

$$\begin{cases} -\Delta u = \lambda \frac{\int_{[0,1]} \alpha e^{\alpha u} \mathcal{P}(d\alpha)}{\iint_{[0,1] \times \Omega} e^{\alpha u} \mathcal{P}(d\alpha) dx} & \text{in } \Omega \\ u = 0 & \text{on } \partial\Omega, \end{cases} \quad (*)_\lambda$$

where $\Omega \subset \mathbb{R}^2$ is a smooth bounded domain, $\lambda > 0$ is a constant and $\mathcal{P} \in \mathcal{M}([0,1])$ is a Borel probability measure. Problem $(*)_\lambda$ was derived by Neri [17] within Onsager's pioneering framework [18], with the aim of including the case of variable vortex intensities. More precisely, in [17] the following mean field equation is derived:

$$\begin{cases} -\Delta v = \frac{\int_{[-1,1]} r e^{-\beta r v} \mathcal{P}(dr)}{\iint_{[-1,1] \times \Omega} e^{-\beta r v} \mathcal{P}(dr) dx} & \text{in } \Omega \\ v = 0 & \text{on } \partial\Omega. \end{cases} \quad (1.1)$$

Here, v is the mean field stream function of an incompressible turbulent Euler flow, the Borel probability measure $\mathcal{P} \in \mathcal{M}([-1,1])$ describes the vortex intensity distribution and $\beta \in \mathbb{R}$ is a constant related to the inverse temperature. The mean field equation (1.1) is derived from the classical Kirchhoff-Routh Hamiltonian for the N -point vortex system:

$$H^N(r_1, \dots, r_N, x_1, \dots, x_N) = \sum_{i \neq j} r_i r_j G(x_i, x_j) + \sum_{i=1}^N r_i^2 H(x_i, x_i),$$

in the limit $N \rightarrow \infty$, under the *stochastic* assumption that the r_i 's are independent identically distributed random variables with distribution \mathcal{P} . In the above formula, for $x, y \in \Omega$, $x \neq y$, $G(x, y)$ denotes the Green's function defined by

$$\begin{cases} -\Delta G(\cdot, y) = \delta_y & \text{in } \Omega \\ G(\cdot, y) = 0 & \text{on } \partial\Omega \end{cases}$$

and $H(x, y)$ denotes the regular part of G , i.e.

$$H(x, y) = G(x, y) + \frac{1}{2\pi} \log |x - y|. \quad (1.2)$$

Setting $u := -\beta v$ and $\lambda = -\beta$, and assuming that

$$\text{supp } \mathcal{P} \subset [0, 1], \quad (1.3)$$

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problem (1.1) takes the form $(*)_\lambda$. We recall that

$$\text{supp } \mathcal{P} := \{\alpha \in [-1, 1] : \mathcal{P}(N) > 0 \text{ for any open neighborhood } N \text{ of } \alpha\}.$$

Assumption (1.3) corresponds to the case of physical interest where all vorticities have the same orientation.

We observe that without loss of generality we may assume

$$1 \in \text{supp } \mathcal{P}. \quad (1.4)$$

Indeed, suppose that $\sup \text{supp } \mathcal{P} = \bar{\alpha} \in (0, 1)$. Then, $(*)_\lambda$ is equivalent to

$$\begin{cases} -\Delta u = \lambda \frac{\int_{[0, \bar{\alpha}]} \alpha e^{\alpha u} \mathcal{P}(d\alpha)}{\iint_{[0, \bar{\alpha}] \times \Omega} e^{\alpha u} \mathcal{P}(d\alpha) dx} & \text{in } \Omega \\ u = 0 & \text{on } \partial\Omega. \end{cases}$$

By the change of variables $\alpha = \alpha' \bar{\alpha}$, $\bar{\mathcal{P}}(A) = \mathcal{P}(\bar{\alpha}A)$ for all Borel sets $A \subset [0, 1]$, and setting $\bar{u} = \bar{\alpha}u$, we find that \bar{u} satisfies

$$\begin{cases} -\Delta \bar{u} = \bar{\alpha}^2 \lambda \frac{\int_{[0, 1]} \alpha' e^{\alpha' \bar{u}} \bar{\mathcal{P}}(d\alpha')}{\iint_{[0, 1] \times \Omega} e^{\alpha' \bar{u}} \bar{\mathcal{P}}(d\alpha') dx} & \text{in } \Omega \\ \bar{u} = 0 & \text{on } \partial\Omega, \end{cases}$$

which is nothing but $(*)_{\bar{\alpha}^2 \lambda}$, with $\bar{\mathcal{P}}$ satisfying (1.4). Henceforth, we always assume (1.4).

When $\mathcal{P}(d\alpha) = \delta_1(d\alpha)$ problem $(*)_\lambda$ reduces to the *standard* mean field problem

$$\begin{cases} -\Delta u = \lambda \frac{e^u}{\int_{\Omega} e^u dx} & \text{in } \Omega \\ u = 0 & \text{on } \partial\Omega, \end{cases}$$

which has been extensively analyzed, see e.g. [13] and the references therein. In the context of turbulence, the case $\mathcal{P}(d\alpha) = \delta_1(d\alpha)$ was developed in [5], see also [3].

Problem $(*)_\lambda$ admits a variational formulation. Indeed, solutions to $(*)_\lambda$ correspond to critical points in $H_0^1(\Omega)$ for the functional

$$J_\lambda(u) = \frac{1}{2} \int_{\Omega} |\nabla u|^2 dx - \lambda \log \left(\iint_{[0, 1] \times \Omega} e^{\alpha u} \mathcal{P}(d\alpha) dx \right). \quad (1.5)$$

Whether or not the optimal value of λ such that J_λ is bounded from below depends on \mathcal{P} was raised as an open question in [26], p. 192, in relation to other apparently similar models for which such a dependence holds true. However, it was noticed in [19] that, in fact, J_λ may be viewed as a perturbation of the standard Moser-Trudinger functional [15, 27] and that, under assumption (1.4), such an optimal value of λ is exactly 8π independently of \mathcal{P} . More precisely, it was already observed in [17] that J_λ is bounded from below on $H_0^1(\Omega)$ if $\lambda \leq 8\pi$. Consequently, the existence of minimizing solutions for $(*)_\lambda$ was obtained in [17] in the subcritical range $\lambda \in (0, 8\pi)$. In [21] the existence of solutions to $(*)_\lambda$ was obtained in the supercritical range $\lambda \in (8\pi, 16\pi)$ under the *non-degeneracy* assumption

$$\mathcal{P}(\{1\}) > 0. \quad (1.6)$$

If (1.6) is satisfied, problem $(*)_\lambda$ may be written in the form $-\Delta u = \rho f(u)$, with $f(t) = e^t + o(e^t)$ as $t \rightarrow +\infty$, $\rho > 0$, and thus it fits into the framework considered in [16, 28]. In particular, if (1.6) is satisfied, the techniques in [16, 28] may be applied to obtain the mass quantization of concentrating solution sequences. On the other hand, the case $\mathcal{P}(\{1\}) = 0$ requires extra care.

Thus, our aim in this note is to complete the existence result in [21] by establishing the existence of solutions to $(*)_\lambda$ in the supercritical regime, *without assuming* (1.6) and for *all* values of λ for which compactness of solution sequences holds.

In order to state our results precisely, we recall that by the Brezis-Merle concentration compactness theory [4], as adapted in [19], an L^∞ -unbounded sequence u_n of solutions to

$(*)_{\lambda_n}$ necessarily concentrates at a finite number of points in Ω , namely

$$\lambda_n \frac{\int_{[0,1]} \alpha e^{\alpha u_n} \mathcal{P}(d\alpha)}{\iint_{[0,1] \times \Omega} e^{\alpha u_n} \mathcal{P}(d\alpha) dx} \xrightarrow{*} \sum_{i=1}^m n_i \delta_{p_i}(dx) + s(x) dx, \quad (1.7)$$

weakly in the sense of $\mathcal{M}(\Omega)$, for some $m \in \mathbb{N}$, $p_i \in \Omega$, $n_i \geq 4\pi$, $i = 1, 2, \dots, m$, and $s \in L^1(\Omega)$. Our first aim is to improve (1.7) by showing that, actually, there holds $n_i = 8\pi$ for all $i = 1, 2, \dots, m$, moreover $s \equiv 0$ and $\lambda_0 \in 8\pi\mathbb{N}$. Namely, we establish the following *mass quantization* result.

Theorem 1.1. *Assume that \mathcal{P} satisfies (1.4). Let $\lambda_n \rightarrow \lambda_0$ and let u_n be a concentrating sequence of solutions to $(*)_{\lambda_n}$. Then, there exist $p_i \in \Omega$, $i = 1, 2, \dots, m$, such that, up to subsequences,*

$$\lambda_n \frac{e^{\alpha u_n} \mathcal{P}(d\alpha)}{\iint_{[0,1] \times \Omega} e^{\alpha u_n} \mathcal{P}(d\alpha) dx} \xrightarrow{*} 8\pi \sum_{i=1}^m \delta_1(d\alpha) \delta_{p_i}(dx), \quad (1.8)$$

weakly in the sense of $\mathcal{M}([0,1] \times \Omega)$. In particular, $\lambda_0 \in 8\pi\mathbb{N}$.

In the *non-degenerate* case (1.6), Theorem 1.1 was established in [21], see also [20] for an alternative proof.

Via Theorem 1.1 and a min-max construction, we shall then obtain the existence result for solutions to $(*)_{\lambda}$. For the existence result we need to assume that Ω is topologically non-trivial, namely that:

$$\Omega \text{ is non-simply connected.} \quad (1.9)$$

Our existence result is the following.

Theorem 1.2. *Assume that \mathcal{P} satisfies (1.4). Assume that Ω satisfies (1.9). Then, for every $\lambda \in \cup_{k \in \mathbb{N}} (8\pi k, 8\pi(k+1))$ there exists a solution to problem $(*)_{\lambda}$.*

We shall obtain the solutions as saddle-type critical points for the Euler-Lagrange functional J_{λ} defined in (1.5), following the variational scheme introduced in [2], see also [8]. It will be clear from the proof that, alternatively, we could follow the variational approach introduced in [10], see also [9].

The article is organized as follows. In Section 2 we recall some known results and we establish some necessary lemmas. In Section 3 we obtain some blow-up results and we prove Theorem 1.1. In Section 4 we set up the variational construction and we prove Theorem 1.2. In the Appendix we show that a suitable rescaling yields a Liouville bubble profile in the limit. This fact, although not needed in the variational construction, provides an intuitive justification to the quantization of the values of λ for which blow-up may occur. In the “degenerate” case $\mathcal{P}(\{1\}) = 0$, the appropriate rescaling parameters depend on \mathcal{P} in a non-trivial way.

Notation. We denote by $C > 0$ a general large constant whose actual value is allowed to vary. We denote by \mathbb{N} the set of positive integers. When the integration variable is clear from the context we omit it. Henceforth, we denote $I := [0, 1]$.

2. PRELIMINARY RESULTS

For the sake of completeness, we collect in this section some preliminary results of various nature which will be used in the sequel.

2.1. Concentration-compactness principle. We recall the Brezis-Merle blow-up theory [4], as adapted to $(*)_{\lambda}$ in [19]. Let us define the sequence of measures $\nu_n \in \mathcal{M}(\Omega)$ by

$$\nu_n(dx) := \lambda_n \frac{\int_{[0,1]} \alpha e^{\alpha u_n(x)} \mathcal{P}(d\alpha)}{\iint_{[0,1] \times \Omega} e^{\alpha u_n(x)} \mathcal{P}(d\alpha) dx} dx.$$

Then, the following alternative holds true.

Lemma 2.1 (Brezis-Merle alternative). *Let u_n be a sequence of solutions to $(*)_{\lambda_n}$ with $\lambda_n \rightarrow \lambda_0$. Then, up to subsequences, exactly one of the following alternatives holds:*

- (i) (*Compactness*) There exists a solution $u_0 \in H_0^1(\Omega)$ to $(*)_{\lambda_0}$ such that $u_n \rightarrow u_0$ in any relevant norm;
- (ii) (*Concentration*) There exists a finite, non-empty blow-up set $\mathcal{S} = \{p_1, \dots, p_m\} \subset \Omega$ such that $u_n \in L_{\text{loc}}^\infty(\Omega \setminus \mathcal{S})$ and

$$\nu_n \xrightarrow{*} \sum_{i=1}^m n_i \delta_{p_i}(dx) + s(x) dx \quad (2.1)$$

for some $n_i \geq 4\pi$, $i = 1, \dots, m$ and for some $s \in L^1(\Omega)$.

Proof. We first observe that, in view of the two-dimensional argument in [12] p. 223, there exists $\varepsilon > 0$, depending only on Ω , such that u_n has no stationary point in an ε -neighborhood of $\partial\Omega$. Consequently, blow-up does not occur on the boundary $\partial\Omega$.

We adapt Theorem 3, p. 1237 in [4] to our case. Let

$$W_n(x) := \lambda_n \frac{\int_{[0,1]} \alpha e^{-(1-\alpha)u_n} \mathcal{P}(d\alpha)}{\iint_{[0,1] \times \Omega} e^{\alpha' u_n} P(d\alpha')}.$$

Then, problem $(*)_{\lambda_n}$ takes the form

$$\begin{cases} -\Delta u_n = W_n(x) e^{u_n} & \text{in } \Omega \\ u_n = 0 & \text{on } \partial\Omega. \end{cases}$$

By the maximum principle, we have $u_n \geq 0$ and hence,

$$0 \leq W_n(x) \leq \frac{\lambda_n}{|\Omega|} \int_{[0,1]} \alpha \mathcal{P}(d\alpha) \quad \text{for all } x \in \Omega.$$

Moreover,

$$\int_{\Omega} W_n(x) e^{u_n} dx \leq \lambda_n. \quad (2.2)$$

Therefore, assumptions (21)–(22) in [4], Theorem 3, are satisfied with the exponent $p = +\infty$. Consequently, it is readily seen that the proof of Theorem 3 in [4] may be adapted in order to prove that either alternative (i) holds true, or there exists a finite set $\mathcal{S} \subset \Omega$ such that, up to subsequences, u_n is bounded in $L_{\text{loc}}^\infty(\Omega \setminus \mathcal{S})$. In the latter case it follows that (2.1) holds true, i.e., alternative (ii) is satisfied. \square

Remark 2.2. In the statement of Theorem 3 in [4], a third assumption (23) is made on the sequence u_n of solutions to $(*)_{\lambda_n}$, namely it is assumed that $\sup_n \int_{\Omega} e^{u_n} < +\infty$. Since in our case we only have the weaker assumption (2.2), we cannot in general directly apply the arguments in [4] to show that $s = 0$ in (2.1). However, if we assume $\mathcal{P}(\{1\}) > 0$, the proof in [4] may be adapted. Indeed, $(*)_{\lambda_n}$ and (2.1) imply that $u_n \rightarrow u_0$ weakly in $W_0^{1,q}(\Omega)$, strongly in $L^q(\Omega)$ for any $1 \leq q < 2$, and a.e., where

$$u_0(x) \geq \sum_{i=1}^m \frac{n_i}{2\pi} \left(\log \frac{1}{|x - p_i|} + H(x, p_i) \right), \quad (2.3)$$

and where H is defined in (1.2). If $\mathcal{P}(\{1\}) > 0$ we may estimate

$$\iint_{[0,1] \times \Omega} e^{\alpha u_n(x)} \mathcal{P}(d\alpha) dx \geq \mathcal{P}(\{1\}) \int_{\Omega} e^{u_n} dx.$$

By Fatou's lemma, (2.3) and recalling that $n_i \geq 4\pi$,

$$\liminf_{n \rightarrow \infty} \int_{\Omega} e^{u_n} \geq \int_{\Omega} e^{u_0} = +\infty$$

and consequently

$$\iint_{[0,1] \times \Omega} e^{\alpha u_n(x)} \mathcal{P}(d\alpha) dx \rightarrow +\infty.$$

This implies $s \equiv 0$.

2.2. Improved Moser-Trudinger inequality. We shall need an improved Moser-Trudinger inequality for the functional (1.5) defined on the bounded domain $\Omega \subset \mathbb{R}^2$.

We recall that the classical Moser-Trudinger sharp inequality [15, 27] states that

$$C_{MT} := \sup \left\{ \int_{\Omega} e^{4\pi u^2} : u \in H_0^1(\Omega), \|\nabla u\|_2 \leq 1 \right\} < +\infty, \quad (2.4)$$

where the constant 4π is best possible. Moreover, the embedding $u \in H_0^1(\Omega) \rightarrow e^u \in L^1(\Omega)$ is compact. For a proof, see, e.g., Theorem 2.46 p. 63 in [1].

In view of the elementary inequality

$$|u| \leq \frac{\|\nabla u\|_2^2}{16\pi} + 4\pi \frac{u^2}{\|\nabla u\|_2^2},$$

we deduce from (2.4) that

$$\log \left(\int_{\Omega} e^{|u|} dx \right) \leq \frac{1}{16\pi} \|\nabla u\|_2^2 + \log(C_{MT}), \quad \text{for any } u \in H_0^1(\Omega). \quad (2.5)$$

In particular, the functional

$$I_{\lambda}(u) = \frac{1}{2} \|\nabla u\|_2^2 - \lambda \log \int_{\Omega} e^u dx$$

is bounded from below for all $\lambda \leq 8\pi$, while it is not difficult to check that

$$\inf_{u \in H_0^1(\Omega)} I_{\lambda}(u) = -\infty \quad (2.6)$$

whenever $\lambda > 8\pi$. Indeed, evaluating the functional I_{λ} on the following adaptation of the Liouville bubbles defined in (5.1) below:

$$u_{\varepsilon}(x) = \begin{cases} \log \frac{(\varepsilon^2 + r_0^2)^2}{(\varepsilon^2 + |x - x_0|^2)^2}, & \text{in } B_{r_0}(x_0) \\ 0, & \text{in } \Omega \setminus B_{r_0}(x_0), \end{cases}$$

yields

$$\int_{\Omega} |\nabla u_{\varepsilon}|^2 dx = 16\pi \log \frac{1}{\varepsilon^2} + O(1), \quad \log \int_{\Omega} e^{u_{\varepsilon}} dx = \log \frac{1}{\varepsilon^2} + O(1)$$

so that

$$I_{\lambda}(u_{\varepsilon}) = (8\pi - \lambda) \log \frac{1}{\varepsilon^2} + O(1) \rightarrow -\infty \quad \text{as } \varepsilon \rightarrow 0.$$

The arguments above imply that if (1.4) is satisfied, then analogous results hold for the Neri's functional J_{λ} . More precisely, we have

Lemma 2.3 (Moser-Trudinger inequality). *Assume that \mathcal{P} satisfies (1.4). Then*

$$\log \iint_{I \times \Omega} e^{\alpha u} \mathcal{P}(d\alpha) dx \leq \frac{1}{16\pi} \|\nabla u\|_2^2 + \log C_{MT} \quad \text{for any } u \in H_0^1(\Omega), \quad (2.7)$$

and the functional $J_{\lambda}(u)$ is bounded from below on $H_0^1(\Omega)$, if and only if $\lambda \leq 8\pi$.

Lemma 2.3 was established in [21] for functions $u \in H^1(M)$ satisfying $\int_M u = 0$, where M is a two-dimensional compact Riemannian manifold. The proof for $u \in H_0^1(\Omega)$ is similar. For the sake of completeness, we outline it below.

On the other hand, in the next Lemma we show that the constant $\frac{1}{16\pi}$ in (2.7) may be lowered if the quantity

$$\frac{\int_I e^{\alpha u(x)} \mathcal{P}(d\alpha)}{\iint_{I \times \Omega} e^{\alpha u} \mathcal{P}(d\alpha) dx},$$

which may be interpreted as the *mass* of u , is suitably distributed. Namely, following ideas of [1, 7], we prove:

Lemma 2.4 (Improved Moser-Trudinger inequality). *Assume that \mathcal{P} satisfies (1.4). Let $d_0 > 0$, $a_0 \in (0, 1/2)$ and for a fixed positive integer ℓ , let $\Omega_1, \dots, \Omega_{\ell+1}$ be subsets of Ω satisfying $\text{dist}(\Omega_i, \Omega_j) \geq d_0$, for all $i \neq j$. Then, for any $\varepsilon > 0$ there exists a constant $K = K(\varepsilon, d_0, a_0, \ell) > 0$ such that if $u \in H_0^1(\Omega)$ satisfies*

$$\frac{\iint_{I \times \Omega_i} e^{\alpha u} \mathcal{P}(d\alpha) dx}{\iint_{I \times \Omega} e^{\alpha u} \mathcal{P}(d\alpha) dx} \geq a_0, \quad i = 1, \dots, \ell + 1, \quad (2.8)$$

then it holds

$$\log \left(\iint_{I \times \Omega} e^{\alpha u} \mathcal{P}(d\alpha) dx \right) \leq \frac{1}{16(\ell + 1)\pi - \varepsilon} \int_{\Omega} |\nabla u|^2 dx + K. \quad (2.9)$$

We begin by outlining the proof of Lemma 2.3.

Proof of Lemma 2.3. The “if” part is immediate and was already used in [17] in order to obtain solutions to $(*)_{\lambda}$ for all $\lambda \in (0, 8\pi)$. Indeed, we have

$$\iint_{I \times \Omega} e^{\alpha u} \mathcal{P}(d\alpha) dx \leq \int_{\Omega} e^{|u|} dx \leq C_{MT} e^{\frac{1}{16\pi} \|\nabla u\|_2^2},$$

for all $u \in H_0^1(\Omega)$. Therefore J_{λ} is bounded below if $\lambda \leq 8\pi$. On the other hand the value 8π is also optimal, provided that $1 \in \text{supp } \mathcal{P}$. In order to show it one needs only to prove that

$$\inf_{u \in H_0^1(\Omega)} J_{\lambda}(u) = -\infty, \quad \text{for any } \lambda > 8\pi. \quad (2.10)$$

Assume (1.4). Since the functional $I_{\lambda}(u)$ is unbounded from below for $\lambda > 8\pi$, then also the functional

$$I_{\lambda}(u)|_{u \geq 0} = \frac{1}{2} \|\nabla u\|_2^2 - \lambda \log \int_{\Omega} e^u dx, \quad u \in H_0^1(\Omega), \quad u \geq 0$$

is unbounded below for $\lambda > 8\pi$. At this point we observe that for every $0 < \delta < 1$ and $u \geq 0$, $u \in H_0^1(\Omega)$, we have:

$$\begin{aligned} J_{\lambda}(u) &= \frac{1}{2} \|\nabla u\|_2^2 - \lambda \log \iint_{I \times \Omega} e^{\alpha u} \mathcal{P}(d\alpha) dx \leq \frac{1}{2} \|\nabla u\|_2^2 - \lambda \log \iint_{[1-\delta, 1] \times \Omega} e^{\alpha u} \mathcal{P}(d\alpha) dx \\ &\leq \frac{1}{2} \|\nabla u\|_2^2 - \lambda \log \int_{\Omega} e^{(1-\delta)u} dx - \lambda \log(\mathcal{P}([1-\delta, 1])) \\ &= \frac{1}{(1-\delta)^2} \left[\frac{1}{2} \|(1-\delta)\nabla u\|_2^2 - \lambda(1-\delta)^2 \log \left(\int_{\Omega} e^{(1-\delta)u} dx \right) \right] - \lambda \log(\mathcal{P}([1-\delta, 1])) \\ &= \frac{1}{(1-\delta)^2} I_{\lambda(1-\delta)^2}((1-\delta)u) - \lambda \log(\mathcal{P}([1-\delta, 1])). \end{aligned}$$

Hence, for $\lambda(1-\delta)^2 > 8\pi$, the right hand side of the last inequality is unbounded from below in view of (2.6), and therefore

$$\inf_{u \in H_0^1(\Omega)} J_{\lambda}(u) = -\infty \quad \text{for any } \lambda > \frac{8\pi}{(1-\delta)^2}.$$

Since $\delta \in (0, 1)$ is arbitrary, (2.10) follows. \square

In order to prove Lemma 2.4, we adapt some ideas contained in [7], Proposition 1.

Proof of Lemma 2.4. Let $g_1, \dots, g_{\ell+1}$ be smooth functions defined on Ω such that $0 \leq g_i \leq 1$, $g_i \equiv 1$ on Ω_i , $|\nabla g_i| \leq c(d_0)$, for $i = 1, \dots, \ell + 1$, and $\text{supp}(g_i) \cap \text{supp}(g_j) = \emptyset$ if $i \neq j$. Up to relabelling, we may assume that

$$\|g_1 \nabla u\|_{L^2(\Omega)} \leq \|g_i \nabla u\|_{L^2(\Omega)} \quad \text{for any } i = 2, \dots, \ell + 1. \quad (2.11)$$

For every $t \in \mathbb{R}$ we denote $t^+ = \max\{0, t\}$. We fix $a > 0$. In view of (2.5), applied to $g_1(|u| - a)^+$, we have that

$$\int_{\Omega} e^{g_1(|u| - a)^+} \leq C_{MT} \exp \left\{ \frac{1}{16\pi} \|\nabla [g_1(|u| - a)^+]\|_{L^2(\Omega)}^2 \right\}. \quad (2.12)$$

Hence, in view of assumption (2.8) and of (2.12), using the elementary inequality $(A+B)^2 \leq (1+\tau)A^2 + c(\tau)B^2$, for any $\tau > 0$, we get

$$\begin{aligned}
\iint_{I \times \Omega} e^{\alpha u} \mathcal{P}(d\alpha) dx &\leq \frac{e^a}{a_0} \iint_{I \times \Omega_1} e^{(\alpha u - a)^+} \mathcal{P}(d\alpha) dx \\
&\leq \frac{e^a}{a_0} \iint_{I \times \Omega} e^{g_1(\alpha u - a)^+} \mathcal{P}(d\alpha) dx \leq \frac{e^a}{a_0} \int_{\Omega} e^{g_1(|u| - a)^+} dx \\
&\leq \frac{C}{a_0} \exp \left\{ \frac{1}{16\pi} \|\nabla [g_1(|u| - a)^+]\|_{L^2(\Omega)}^2 + a \right\} \\
&\leq \frac{C}{a_0} \exp \left\{ \frac{1}{16\pi} \left[(1+\tau) \|g_1 \nabla u\|_{L^2(\Omega)}^2 + c(\tau) \|(|u| - a)^+ \nabla g_1\|_{L^2(\Omega)}^2 \right] + a \right\} \\
&\stackrel{(2.11)}{\leq} \frac{C}{a_0} \exp \left\{ \frac{1}{16\pi(\ell+1)} \left[(1+\tau) \sum_{i=1}^{\ell+1} \|g_i \nabla u\|_{L^2(\Omega)}^2 + c(\tau, d_0) \|(|u| - a)^+\|_{L^2(\Omega)}^2 \right] + a \right\} \\
&\leq \frac{C}{a_0} \exp \left\{ \frac{1}{16\pi(\ell+1)} \left[(1+\tau) \|\nabla u\|_{L^2(\Omega)}^2 + c(\tau, d_0) \|(|u| - a)^+\|_{L^2(\Omega)}^2 \right] + a \right\},
\end{aligned}$$

where we used $\text{supp}(g_i) \cap \text{supp}(g_j) = \emptyset$ to derive the last inequality and where $C = C_{MT}$.

For a given real number $\eta \in (0, |\Omega|)$, let a be such that $\text{meas}(\{x \in \Omega : |u(x)| \geq a\}) = \eta$. Then, by the Hölder and Sobolev inequalities we have

$$\|(|u| - a)^+\|_{L^2(\Omega)}^2 = \eta^{1/2} \left(\int_{\{x \in \Omega : |u| > a\}} (|u| - a)^4 \right)^{1/2} \leq \eta^{1/2} C \|\nabla u\|_2^2.$$

Using the Schwarz and Poincaré inequalities, we finally derive

$$a\eta \leq \int_{\{x \in \Omega : |u| \geq a\}} |u| \leq |\Omega|^{1/2} \left(\int_{\Omega} |u|^2 \right)^{1/2} \leq C \|\nabla u\|_2,$$

and therefore, for any small $\delta > 0$,

$$a \leq \frac{\delta}{2} \|\nabla u\|_2^2 + \frac{C^2}{2\delta\eta^2}.$$

In conclusion, we have derived that

$$\iint_{I \times \Omega} e^{\alpha u} \mathcal{P}(d\alpha) dx \leq \frac{C}{a_0} \exp \left\{ \frac{1}{16\pi(\ell+1)} \left(1 + \tau + c(\tau, d_0) \eta^{\frac{1}{2}} C + \frac{\delta}{2} \right) \|\nabla u\|_{L^2(\Omega)}^2 + \frac{C^2}{2\delta\eta^2} \right\}.$$

Let $\tilde{\varepsilon} = \frac{\varepsilon}{16\pi(\ell+1)-\varepsilon}$. Fixing $\tau < \frac{\tilde{\varepsilon}}{3}$, η such that $c(\tau, d_0) \eta^{\frac{1}{2}} C < \frac{\tilde{\varepsilon}}{3}$ and δ such that $\frac{\delta}{2} < \frac{\tilde{\varepsilon}}{3}$, the asserted improved Moser-Trudinger inequality (2.9) is completely established. \square

Using Lemma 2.4 we can characterize the limiting behavior of sequences of measures on Ω of the form

$$\frac{\int_I e^{\alpha u_n} \mathcal{P}(d\alpha) dx}{\iint_{I \times \Omega} e^{\alpha u_n} \mathcal{P}(d\alpha) dx}, \quad (2.13)$$

where the functions u_n are such that the functional J_λ attains arbitrarily large negative values. Such a characterization will be used in an essential way in the variational scheme, in particular in the proof of Proposition 4.3 below.

Lemma 2.5 (Concentration property). *Assume that \mathcal{P} satisfies (1.4). Let $\lambda \in (8k\pi, 8(k+1)\pi)$, $k \geq 1$, and let $\{u_n\} \subset H_0^1(\Omega)$ be a sequence of functions satisfying $J_\lambda(u_n) \rightarrow -\infty$. For any $\varepsilon > 0$ and for any $r > 0$, there exists a subsequence $\{u_{n_j}\}$ (depending only on ε and r) and ℓ points, $\ell \in \{1, \dots, k\}$, $p_1, \dots, p_\ell \in \overline{\Omega}$ (which do not depend on j) such that*

$$\frac{\iint_{I \times (\Omega \setminus \bigcup_{i=1}^{\ell} B_r(p_i))} e^{\alpha u_{n_j}} \mathcal{P}(d\alpha) dx}{\iint_{I \times \Omega} e^{\alpha u_{n_j}} \mathcal{P}(d\alpha) dx} < \varepsilon \quad \text{for any } i \in \{1, \dots, \ell\} \text{ and for any } j \quad (2.14)$$

and

$$\lim_{j \rightarrow +\infty} \left(\frac{\iint_{I \times (B_r(p_i) \setminus \bigcup_{h=1}^{i-1} B_r(p_h))} e^{\alpha u_{n_j}} \mathcal{P}(d\alpha) dx}{\iint_{I \times \Omega} e^{\alpha u_{n_j}} \mathcal{P}(d\alpha) dx} \right) = \beta_i > 0 \quad \text{for any } i \in \{1, \dots, \ell\}, \quad (2.15)$$

with $\sum_{i=1}^{\ell} \beta_i = 1$.

Proof. The proof is a direct consequence of some general concentration properties of L^1 -functions, as obtained in Lemma 3.3 in [14] and Lemma 2.4 in [2], applied to the functions (2.13). \square

2.3. Properties of some Vandermonde-type maps. Finally, we collect some results from [2] concerning Vandermonde-type maps. Such properties will be needed in order to perform the variational scheme in the proof of Theorem 1.2, and in particular to prove Proposition 4.3 below.

Henceforth, for $k \in \mathbb{N}$, we denote by \underline{z}_k the vector $\underline{z}_k = (z_1, \dots, z_k) \in \mathbb{C}^k$. In particular we denote by $\underline{0}_k \in \mathbb{C}^k$ the null vector, i.e., the vector whose entries are all equal to 0 $\in \mathbb{C}$.

Let D_k be the open unit ball of \mathbb{C}^k , namely

$$D_k = \{\underline{z}_k \in \mathbb{C}^k \mid |z_1|^2 + \dots + |z_k|^2 < 1\}.$$

Let $\Phi_k : \mathbb{C}^k \mapsto \mathbb{C}^k$ be the continuous map defined as

$$\Phi_k(\underline{z}_k) := \begin{pmatrix} z_1 |z_1| & + & z_2 |z_2| & + & \dots & + & z_k |z_k| \\ z_1^2 & + & z_2^2 & + & \dots & + & z_k^2 \\ \vdots & & \vdots & & \vdots & & \vdots \\ z_1^2 \left(\frac{z_1}{|z_1|}\right)^{j-2} & + & z_2^2 \left(\frac{z_2}{|z_2|}\right)^{j-2} & + & \dots & + & z_k^2 \left(\frac{z_k}{|z_k|}\right)^{j-2} \\ \vdots & & \vdots & & \vdots & & \vdots \\ z_1^2 \left(\frac{z_1}{|z_1|}\right)^{k-2} & + & z_2^2 \left(\frac{z_2}{|z_2|}\right)^{k-2} & + & \dots & + & z_k^2 \left(\frac{z_k}{|z_k|}\right)^{k-2} \end{pmatrix}. \quad (2.16)$$

In [2], Lemma 4.1, the degree of Φ_k was considered and the following was established.

Lemma 2.6. *If $k \in \mathbb{N}$, then*

$$\deg(\Phi_k, \partial D_k, \underline{0}_k) \neq 0.$$

Next we recall another useful result obtained in [2], Lemma 3.3.

Lemma 2.7. *Let $\ell \in \mathbb{N}$ and $\underline{\beta}_\ell \in \mathbb{R}^\ell$ such that $\beta_i > 0$ for any $i \in \{1, \dots, \ell\}$. Suppose that $\underline{z}_\ell \in \mathbb{C}^\ell$ is a solution to*

$$\begin{cases} \beta_1 z_1 + \beta_2 z_2 + \dots + \beta_\ell z_\ell = y_1 \\ \beta_1 z_1^2 + \beta_2 z_2^2 + \dots + \beta_\ell z_\ell^2 = y_2 \\ \dots \dots \dots \\ \beta_1 z_1^\ell + \beta_2 z_2^\ell + \dots + \beta_\ell z_\ell^\ell = y_\ell \end{cases}$$

where $\underline{y}_\ell \in \mathbb{C}^\ell$. Then $\underline{z}_\ell \rightarrow \underline{0}_\ell$ as $\underline{y}_\ell \rightarrow \underline{0}_\ell$.

3. BLOW-UP ANALYSIS AND PROOF OF THEOREM 1.1

Let u_n be a concentrating sequence of solutions to $(*)_{\lambda_n}$. In order to prove Theorem 1.1, we define the sequence of measures $\mu_n \in \mathcal{M}([0, 1] \times \Omega)$ on the product space $[0, 1] \times \Omega$:

$$\mu_n(d\alpha dx) := \lambda_n \frac{e^{\alpha u_n(x)}}{\iint_{[0, 1] \times \Omega} e^{\alpha u_n(x)} \mathcal{P}(d\alpha) dx} \mathcal{P}(d\alpha) dx.$$

Clearly, $\mu_n([0, 1] \times \Omega) = \lambda_n$, therefore there exists a measure $\mu \in \mathcal{M}([0, 1] \times \Omega)$ such that, up to subsequences, $\mu_n \xrightarrow{*} \mu$ weakly in $\mathcal{M}([0, 1] \times \Omega)$. In view of the Brezis-Merle theory [4], as adapted in Lemma 2.1, there exists a finite set $\mathcal{S} = \{p_1, p_2, \dots, p_m\} \subset \Omega$ such that the singular part of μ is supported on $[0, 1] \times \mathcal{S}$. It follows that there exist $\zeta_i \in \mathcal{M}([0, 1])$, $i = 1, \dots, m$, and $r \in L^1([0, 1] \times \Omega)$ such that the limit measure μ is of the form:

$$\mu(d\alpha dx) = \sum_{i=1}^m \zeta_i(d\alpha) \delta_{p_i}(dx) + r(\alpha, x) \mathcal{P}(d\alpha) dx. \quad (3.1)$$

With this notation, the main ingredient in the proof of Theorem 1.1 is the following result.

Proposition 3.1. *Assume (1.4). Let u_n be a sequence of solutions to $(*)_{\lambda_n}$ with $\lambda_n \rightarrow \lambda_0$ and suppose that $\mathcal{S} \neq \emptyset$. Then,*

- (i) $\zeta_i(d\alpha) = 8\pi\delta_1(d\alpha)$, for all $i = 1, 2, \dots, m$.
- (ii) $r \equiv 0$.

In order to establish Proposition 3.1, we first show that the measures $\zeta_i(d\alpha)$ are concentrated at 1 (see Lemma 3.2 below), namely that $\zeta_i(d\alpha) = n_i\delta_1(d\alpha)$ for some $n_i > 0$. Next we provide a quadratic identity for the blow up measures ζ_i (see Lemma 3.3 below), which will involve that $n_i = 8\pi$ for any $i = 1, \dots, m$. In turn, even without the additional assumption $\mathcal{P}(\{1\}) > 0$, the argument outlined in Remark 2.2, to show that $s \equiv 0$ (where s is defined in (2.1)), will allow us to conclude that $r \equiv 0$.

Let us state and prove two preliminary lemmas.

Fix $\varepsilon > 0$. For every $\alpha \in [0, 1 - \varepsilon)$ we define for $x \in \Omega$

$$f_\alpha(x) := \frac{\lambda e^{\alpha u}}{\iint_{I \times \Omega} e^{\alpha' u} \mathcal{P}(d\alpha') dx}.$$

Lemma 3.2. *Assume (1.4). For every $\alpha \in [0, 1 - \varepsilon)$ the following estimate holds:*

$$\int_{\Omega} f_\alpha^{(1-\varepsilon)/\alpha} dx \leq \left(\frac{\lambda}{\mathcal{P}([1-\varepsilon, 1])|\Omega|} \right)^{(1-\varepsilon)/\alpha} |\Omega|.$$

Proof. By definition,

$$\int_{\Omega} f_\alpha^{(1-\varepsilon)/\alpha} dx = \left(\frac{\lambda}{\iint_{I \times \Omega} e^{\alpha' u} \mathcal{P}(d\alpha') dx} \right)^{(1-\varepsilon)/\alpha} \int_{\Omega} e^{(1-\varepsilon)u} dx. \quad (3.2)$$

We observe that, since $u \geq 0$ by the maximum principle, we have $1 \leq e^{(1-\varepsilon)u} \leq e^{\alpha' u}$, and therefore

$$\int_{\Omega} e^{(1-\varepsilon)u} dx \leq \frac{1}{\mathcal{P}([1-\varepsilon, 1])} \iint_{[1-\varepsilon, 1] \times \Omega} e^{\alpha' u} \mathcal{P}(d\alpha') dx. \quad (3.3)$$

We also obtain that

$$\iint_{[1-\varepsilon, 1] \times \Omega} e^{\alpha' u} \mathcal{P}(d\alpha') dx \geq \mathcal{P}([1-\varepsilon, 1])|\Omega|$$

and consequently, recalling that $(1-\varepsilon)/\alpha > 1$,

$$\begin{aligned} \left(\frac{1}{\mathcal{P}([1-\varepsilon, 1])|\Omega|} \iint_{[1-\varepsilon, 1] \times \Omega} e^{\alpha' u} \mathcal{P}(d\alpha') dx \right)^{(1-\varepsilon)/\alpha} \\ \geq \frac{1}{\mathcal{P}([1-\varepsilon, 1])|\Omega|} \iint_{[1-\varepsilon, 1] \times \Omega} e^{\alpha' u} \mathcal{P}(d\alpha') dx. \end{aligned}$$

In turn, we derive that

$$\begin{aligned} \left(\frac{1}{\iint_{[0, 1] \times \Omega} e^{\alpha' u} \mathcal{P}(d\alpha') dx} \right)^{(1-\varepsilon)/\alpha} &\leq \left(\frac{1}{\iint_{[1-\varepsilon, 1] \times \Omega} e^{\alpha' u} \mathcal{P}(d\alpha') dx} \right)^{(1-\varepsilon)/\alpha} \\ &\leq \frac{(\mathcal{P}([1-\varepsilon, 1])|\Omega|)^{1-\frac{1-\varepsilon}{\alpha}}}{\iint_{[1-\varepsilon, 1] \times \Omega} e^{\alpha' u} \mathcal{P}(d\alpha') dx}. \end{aligned} \quad (3.4)$$

Inserting (3.3) and (3.4) into (3.2), we obtain

$$\begin{aligned} \int_{\Omega} f_\alpha^{(1-\varepsilon)/\alpha} dx &\leq \lambda^{(1-\varepsilon)/\alpha} \frac{(\mathcal{P}([1-\varepsilon, 1])|\Omega|)^{1-\frac{1-\varepsilon}{\alpha}}}{\iint_{[1-\varepsilon, 1] \times \Omega} e^{\alpha' u} \mathcal{P}(d\alpha') dx} \cdot \frac{1}{\mathcal{P}([1-\varepsilon, 1])} \iint_{[1-\varepsilon, 1] \times \Omega} e^{\alpha' u} \mathcal{P}(d\alpha') dx \\ &= \left(\frac{\lambda}{\mathcal{P}([1-\varepsilon, 1])|\Omega|} \right)^{(1-\varepsilon)/\alpha} |\Omega|, \end{aligned}$$

as asserted. \square

Lemma 3.3. *For every $i = 1, \dots, m$ the following identity holds:*

$$8\pi \int_{[0,1]} \zeta_i(d\alpha) = \left[\int_{[0,1]} \alpha \zeta_i(d\alpha) \right]^2. \quad (3.5)$$

Proof. See [19], Theorem 2.2, where (3.5) is derived in a more general context by using a symmetry argument introduced in [23]. Alternatively, (3.5) may be derived from the classical Pohozaev identity, see, e.g., [13]. \square

Now we can prove Proposition 3.1.

Proof of Proposition 3.1. Proof of (i). In view of Lemma 3.2, for any $\varepsilon > 0$ we have

$$\begin{aligned} & \iint_{[0,1-\varepsilon] \times \Omega} \left(\frac{\lambda e^{\alpha u_n}}{\iint_{I \times \Omega} e^{\alpha' u_n} \mathcal{P}(d\alpha') dx} \right)^{(1-\varepsilon/2)/(1-\varepsilon)} \mathcal{P}(d\alpha) dx \\ &= \int_{[0,1-\varepsilon]} \mathcal{P}(d\alpha) \int_{\Omega} f_{\alpha}^{(1-\varepsilon/2)/(1-\varepsilon)} dx \\ &\leq \int_{[0,1-\varepsilon]} \left(\int_{\Omega} f_{\alpha}^{(1-\varepsilon/2)/\alpha} dx \right)^{\alpha/(1-\varepsilon)} |\Omega|^{1-\alpha/(1-\varepsilon)} \mathcal{P}(d\alpha) \\ &\leq \left\{ \left(\frac{\lambda}{\mathcal{P}([1-\varepsilon, 1])|\Omega|} \right)^{(1-\varepsilon/2)/\alpha} |\Omega| \right\}^{\alpha/(1-\varepsilon)} |\Omega|^{1-\alpha/(1-\varepsilon)} \\ &\leq \left(\frac{\lambda}{\mathcal{P}([1-\varepsilon, 1])|\Omega|} \right)^{(1-\varepsilon/2)/(1-\varepsilon)} |\Omega|. \end{aligned}$$

It follows that the sequence of functions

$$\mu_n(\alpha, x) = \frac{\lambda e^{\alpha u_n}}{\iint_{I \times \Omega} e^{\alpha' u_n} \mathcal{P}(d\alpha') dx}$$

is uniformly bounded in $L^{(1-\varepsilon/2)/(1-\varepsilon)}([0, 1-\varepsilon] \times \Omega)$. Therefore, for all $i = 1, \dots, m$ we have $\mu(d\alpha dx)|_{[0,1-\varepsilon] \times \Omega} \equiv 0$ and $\text{supp}(\zeta_i) \subset [1-\varepsilon, 1]$ for any $\varepsilon > 0$. This implies that $\zeta_i(d\alpha) = n_i \delta_1(d\alpha)$ for some $n_i > 0$. In turn, from Lemma 3.3 we find $8\pi n_i = (n_i^2)$ and therefore $n_i = 8\pi$.

Proof of (ii). We have, for any sufficiently small $\varepsilon > 0$:

$$\iint_{[0,1] \times \Omega} e^{\alpha u_n} \mathcal{P}(d\alpha) dx \geq \mathcal{P}([1-\varepsilon, 1]) \int_{\Omega} e^{(1-\varepsilon)u_n} dx.$$

On the other hand, up to subsequences, $u_n \rightarrow u_0$ in $W_0^{1,q}(\Omega)$ for any $q \in [1, 2)$, where in view of (2.3) and Part (i) there holds

$$u_0(x) \geq 4 \sum_{i=1}^m \left(\log \frac{1}{|x - p_i|} + H(x, p_i) \right), \quad \text{in } \Omega \setminus \mathcal{S}.$$

In particular,

$$e^{(1-\varepsilon)u_0(x)} \geq \prod_{i=1}^m \frac{c_0}{|x - p_i|^{4(1-\varepsilon)}},$$

and therefore $\int_{\Omega} e^{(1-\varepsilon)u_0} = +\infty$. Hence, by Fatou's lemma we conclude that

$$\iint_{[0,1] \times \Omega} e^{\alpha u_n} \mathcal{P}(d\alpha) dx \rightarrow +\infty \quad (3.6)$$

along a blow-up sequence. Since by Lemma 2.1 $u_n \in L_{\text{loc}}^{\infty}(\Omega \setminus \mathcal{S})$, (3.6) implies $r \equiv 0$ in $[0, 1] \times \Omega$. \square

Proof of Theorem 1.1. In view of Proposition 3.1, the limit (1.8) holds true. We are only left to check the quantization property $\lambda_0 \in 8\pi\mathbb{N}$. This fact readily follows from (1.8). \square

4. THE MIN-MAX SCHEME AND THE PROOF OF THEOREM 1.2

Finally, in this section we complete the proof of Theorem 1.2. In view of assumption (1.9) there exists a simple, closed, smooth, non-contractible curve $\Gamma \subset \Omega$. Moreover, by the Jordan-Schoenflies Theorem [22], there exists a diffeomorphism

$$\chi : \mathbb{R}^2 \rightarrow \mathbb{C}$$

such that $\chi(\Gamma) = \partial B_1(0)$ and such that the bounded component of $\mathbb{R}^2 \setminus \Gamma$ is mapped in $B_1(0)$. Clearly, there exists a point $z_0 \in B_1(0) \subset \mathbb{C}$ and a radius $\rho > 0$ such that $B_{2\rho}(z_0) \cap \chi(\Omega) = \emptyset$. Without loss of generality we may assume that $z_0 = 0$, so that in conclusion we have

$$\partial B_1(0) = \chi(\Gamma) \subset \chi(\Omega), \quad \text{while} \quad B_{2\rho}(0) \cap \chi(\Omega) = \emptyset. \quad (4.1)$$

Via χ , we may also define a simple, regular parametrization of Γ :

$$\gamma : [0, 2\pi) \rightarrow \Gamma, \quad \gamma(\theta) = \chi^{-1}(e^{i\theta}).$$

For $u \in H_0^1(\Omega)$ and $j \in \mathbb{N}$, let $m_j : H_0^1(\Omega) \rightarrow \mathbb{C}$ be defined by

$$m_j(u) = \frac{\iint_{I \times \Omega} (\chi(x))^j e^{\alpha u(x)} \mathcal{P}(d\alpha) dx}{\iint_{I \times \Omega} e^{\alpha' u(x)} \mathcal{P}(d\alpha') dx} = \int_{\Omega} (\chi(x))^j d\mu(u),$$

where

$$d\mu(u) = \frac{\int_I e^{\alpha u(x)} \mathcal{P}(d\alpha)}{\iint_{I \times \Omega} e^{\alpha' u(x)} \mathcal{P}(d\alpha') dx}.$$

For $k \in \mathbb{N}$, let $m : H_0^1(\Omega) \rightarrow \mathbb{C}^k$ be the vectorial map

$$m(u) = (m_1(u), m_2(u), \dots, m_k(u)).$$

We now define, for $k \in \mathbb{N}$, the class of functions which will be used in the min-max argument:

$$\mathcal{F}_\lambda = \left\{ h \in \mathcal{C}(D_k, H_0^1(\Omega)) \mid \begin{array}{l} (i) \ J_\lambda(h(\underline{z}_k)) \rightarrow -\infty \text{ as } \underline{z}_k \rightarrow \partial D_k, \\ (ii) \ m \circ h \text{ can be extended continuously to } \overline{D_k} \\ (iii) \ m \circ h : \partial D_k \rightarrow \mathbb{C}^k \text{ has non zero degree} \end{array} \right\}. \quad (4.2)$$

Proposition 4.1. *Assume (1.4) and let $k \in \mathbb{N}$. For any $\lambda \in (8k\pi, 8(k+1)\pi)$ the set \mathcal{F}_λ is non-empty.*

In order to prove Proposition 4.1, we define a suitable test function. Let $\varepsilon_0 > 0$ such that $B_{\varepsilon_0}(\gamma(\theta)) \subset \Omega$ for any $\theta \in [0, 2\pi)$ and for $(r, \theta) \in [0, 1) \times [0, 2\pi)$ let

$$v_{r,\theta}(x) = \begin{cases} 0 & \text{if } x \in \Omega \setminus B_{\varepsilon_0}(\gamma(\theta)) \\ 4 \log\left(\frac{\varepsilon_0}{|x - \gamma(\theta)|}\right) & \text{if } x \in B_{\varepsilon_0}(\gamma(\theta)) \setminus B_{\varepsilon_0(1-r)}(\gamma(\theta)) \\ 4 \log\left(\frac{1}{1-r}\right) & \text{if } x \in B_{\varepsilon_0(1-r)}(\gamma(\theta)). \end{cases}$$

Let us consider, for $k \in \mathbb{N}$, the following family of probability measures, known in the literature as the set of formal barycenters of Γ of order k :

$$\Gamma_k := \left\{ \sum_{i=1}^k t_i \delta_{\gamma(\theta_i)} : t_i \in [0, 1], \sum_{i=1}^k t_i = 1, \theta_i \in [0, 2\pi) \right\}.$$

Let us fix $\tilde{\alpha} \in I$ satisfying

$$\tilde{\alpha} \in \left(\frac{3}{4}, 1\right) \quad \text{and} \quad (2\tilde{\alpha} - 1)\lambda > 8k\pi. \quad (4.3)$$

Then, given $\sigma \in \Gamma_k$, $\sigma = \sum_{i=1}^k t_i \delta_{\gamma(\theta_i)}$ and $r \in [0, 1)$, we define the function $u_{r,\sigma} \in H_0^1(\Omega)$ by

$$u_{r,\sigma}(x) = \frac{1}{\tilde{\alpha}} \log \left(\sum_{i=1}^k t_i e^{\tilde{\alpha} v_{r,\theta_i}} \right). \quad (4.4)$$

It is readily checked that $u_{r,\sigma}$ depends continuously on $r \in [0, 1)$ and $\sigma \in \Gamma_k$.

Lemma 4.2. Assume (1.4). For $\lambda \in (8k\pi, 8(k+1)\pi)$, $r \in [0, 1)$ and $\sigma \in \Gamma_k$, then

$$J_\lambda(u_{r,\sigma}) \longrightarrow -\infty \quad \text{as } r \rightarrow 1, \quad \text{uniformly for } \sigma \in \Gamma_k. \quad (4.5)$$

and

$$d\mu(u_{r,\sigma}) = \frac{\int_I e^{\alpha u_{r,\sigma}} \mathcal{P}(d\alpha)}{\iint_{I \times \Omega} e^{\alpha u_{r,\sigma}} \mathcal{P}(d\alpha) dx} \xrightarrow{*} \sigma \quad \text{as } r \rightarrow 1, \quad \text{uniformly for } \sigma \in \Gamma_k, \quad (4.6)$$

where the function $u_{r,\sigma}$ is defined in (4.4) with $\tilde{\alpha}$ satisfying (4.3).

Proof. Recalling that

$$J_\lambda(u_{r,\sigma}) = \frac{1}{2} \int_\Omega |\nabla u_{r,\sigma}|^2 dx - \lambda \log \left(\iint_{I \times \Omega} e^{\alpha u_{r,\sigma}} \mathcal{P}(d\alpha) dx \right),$$

property (4.5) will follow from the following two estimates:

$$\int_\Omega |\nabla u_{r,\sigma}|^2 dx \leq 32k\pi \log \frac{1}{1-r} + O(1), \quad (4.7)$$

$$\log \left(\iint_{I \times \Omega} e^{\alpha u_{r,\sigma}} \mathcal{P}(d\alpha) dx \right) = (4\tilde{\alpha} - 2) \log \frac{1}{1-r} + \log(\mathcal{P}([\tilde{\alpha}, 1])) + O(1). \quad (4.8)$$

We note that $\mathcal{P}([\tilde{\alpha}, 1]) > 0$ in view of assumption 1.4. Indeed, from (4.7)–(4.8) and recalling (4.3), it follows that

$$J_\lambda(u_{r,\sigma}) \leq 2(8k\pi - (2\tilde{\alpha} - 1)\lambda) \log \frac{1}{1-r} - \lambda \log(\mathcal{P}([\tilde{\alpha}, 1])) + O(1) \rightarrow -\infty \quad \text{as } r \rightarrow 1.$$

Proof of (4.7). By definition of v_{r,θ_i} , $u_{r,\sigma}$ as in (4)–(4.4), we have

$$\nabla u_{r,\sigma}(x) = \frac{\sum_{i=1}^k t_i e^{\tilde{\alpha} v_{r,\theta_i}(x)} \nabla v_{r,\theta_i}(x)}{\sum_{i=1}^k t_i e^{\tilde{\alpha} v_{r,\theta_i}(x)}},$$

and

$$|\nabla v_{r,\theta_i}| = \begin{cases} \frac{4}{|x - \gamma(\theta_i)|} & x \in B_{\varepsilon_0}(\gamma(\theta_i)) \setminus B_{\varepsilon_0(1-r)}(\gamma(\theta_i)) \\ 0 & \text{otherwise.} \end{cases}$$

Therefore, if $x \in \Omega \setminus \cup_{j=1}^k B_{\varepsilon_0(1-r)}(\gamma(\theta_j))$, then

$$|\nabla u_{r,\sigma}(x)| \leq \frac{\sum_{i=1}^k t_i e^{\tilde{\alpha} v_{r,\theta_i}(x)} \frac{4}{|x - \gamma(\theta_i)|}}{\sum_{i=1}^k t_i e^{\tilde{\alpha} v_{r,\theta_i}(x)}} \leq \frac{4}{\min_{i=1,2,\dots,k} |x - \gamma(\theta_i)|}.$$

From the estimate above we also deduce that, for any $x \in \Omega$,

$$|\nabla u_{r,\sigma}(x)| \leq \frac{4}{\varepsilon_0(1-r)}.$$

Then setting $A_i = \{y \in \Omega : |y - \gamma(\theta_i)| = \min_j |y - \gamma(\theta_j)|\}$, we have

$$\begin{aligned} \int_\Omega |\nabla u_{r,\sigma}(x)|^2 dx &\stackrel{(4)}{\leq} \sum_{i=1}^k \int_{B_{\varepsilon_0(1-r)}(\gamma(\theta_i))} |\nabla u_{r,\sigma}(x)|^2 dx + \int_{\Omega \setminus \cup_{j=1}^k B_{\varepsilon_0(1-r)}(\gamma(\theta_j))} \frac{16}{(\min_i |x - \gamma(\theta_i)|)^2} dx \\ &\stackrel{(4)}{\leq} O(1) + \sum_{i=1}^k \int_{A_i \setminus \cup_{j=1}^k B_{\varepsilon_0(1-r)}(\gamma(\theta_j))} \frac{16}{|x - \gamma(\theta_i)|^2} dx \\ &\leq O(1) + \sum_{i=1}^k \int_{A_i \setminus B_{\varepsilon_0(1-r)}(\gamma(\theta_i))} \frac{16}{|x - \gamma(\theta_i)|^2} dx \\ &\leq O(1) + 16k \int_{B_{\text{diam}(\Omega)}(\gamma(\theta_i)) \setminus B_{\varepsilon_0(1-r)}(\gamma(\theta_i))} \frac{dx}{|x - \gamma(\theta_i)|^2} \\ &\leq O(1) + 32k\pi \log \frac{1}{1-r}, \end{aligned}$$

so that (4.7) is proved.

Proof of (4.8). It is readily checked that

$$\begin{aligned} \log \left(\iint_{I \times \Omega} e^{\alpha u_{r,\sigma}} \mathcal{P}(d\alpha) dx \right) &\geq \log \left(\iint_{[\tilde{\alpha}, 1] \times \Omega} e^{\tilde{\alpha} u_{r,\sigma}} \mathcal{P}(d\alpha) dx \right) \\ &= \log \left(\int_{\Omega} e^{\tilde{\alpha} u_{r,\sigma}} dx \right) + \log(\mathcal{P}([\tilde{\alpha}, 1])). \end{aligned} \quad (4.9)$$

Then, recalling that in view of (4.3) we have $\tilde{\alpha} > \frac{1}{2}$,

$$\begin{aligned} \int_{\Omega} e^{\tilde{\alpha} u_{r,\sigma}} dx &= \int_{\Omega} \sum_{i=1}^k t_i e^{\tilde{\alpha} v_{r,\theta_i}} dx \\ &= \sum_{i=1}^k t_i \left[\int_{B_{\varepsilon_0(1-r)}(\gamma(\theta_i))} \frac{dx}{(1-r)^{4\tilde{\alpha}}} + \int_{B_{\varepsilon_0}(\gamma(\theta_i)) \setminus B_{\varepsilon_0(1-r)}(\gamma(\theta_i))} \left(\frac{\varepsilon_0}{|x - \gamma(\theta_i)|} \right)^{4\tilde{\alpha}} dx \right. \\ &\quad \left. + \int_{\Omega \setminus B_{\varepsilon_0}(\gamma(\theta_i))} dx \right] \\ &= \sum_{i=1}^k t_i \left[\pi \varepsilon_0^2 (1-r)^{2-4\tilde{\alpha}} + 2\pi \varepsilon_0^{4\tilde{\alpha}} \int_{\varepsilon_0(1-r)}^{\varepsilon_0} \frac{d\rho}{\rho^{4\tilde{\alpha}-1}} + |\Omega| - \pi \varepsilon_0^2 \right] \\ &= \frac{C}{(1-r)^{4\tilde{\alpha}-2}} + O(1). \end{aligned} \quad (4.10)$$

Finally, combining (4.9) and (4.10) we obtain (4.8). Hence, (4.5) is completely established.

Proof of (4.6). Let $\varepsilon_0 > 0$ be such that $B_{\varepsilon_0}(\gamma(\theta)) \subset \Omega$ for any $\theta \in [0, 2\pi)$ and let $\sigma = \sum_{i=1}^k t_i \delta_{\gamma(\theta_i)}$. Without loss of generality we may assume that there exists $m = m(\sigma) \leq k$ such that $t_i > 0$ for any $i = 1, \dots, m$ and $t_i = 0$ for $i > m$.

In order to prove (4.6) it suffices to show that for every $\varepsilon \in (0, \varepsilon_0)$

$$\lim_{r \rightarrow 1} \int_{\cup_{i=1}^m B_{\varepsilon}(\gamma(\theta_i))} d\mu(u_{r,\sigma}) = 1 \quad \text{uniformly with respect to } \sigma \in \Gamma_k. \quad (4.11)$$

Let us fix $\varepsilon \in (0, \varepsilon_0)$, let $\delta = \delta(\varepsilon) \in (0, 1)$ such that $B_{\delta}(0) \subset \varphi_{\theta}(B_{\varepsilon}(\gamma(\theta)))$, where $\varphi_{\theta}(x) = \frac{x - \gamma(\theta)}{\varepsilon_0}$.

We write

$$\int_{\cup_{i=1}^m B_{\varepsilon}(\gamma(\theta_i))} d\mu(u_{r,\sigma}) = \frac{A+B}{A+C} \quad (4.12)$$

where

$$A = \iint_{I \times \cup_{i=1}^m \varphi_{\theta_i}^{-1}(B_{\delta}(0))} e^{\alpha u_{r,\sigma}} \mathcal{P}(d\alpha) dx,$$

and

$$\begin{aligned} B &= \iint_{I \times (\cup_{i=1}^m (B_{\varepsilon}(\gamma(\theta_i)) \setminus \varphi_{\theta_i}^{-1}(B_{\delta}(0))))} e^{\alpha u_{r,\sigma}} \mathcal{P}(d\alpha) dx, \\ C &= \iint_{I \times (\Omega \setminus \cup_{i=1}^m \varphi_{\theta_i}^{-1}(B_{\delta}(0)))} e^{\alpha u_{r,\sigma}} \mathcal{P}(d\alpha) dx. \end{aligned}$$

We claim that, as $r \rightarrow 1$,

$$A \rightarrow +\infty, \quad B = O(1), \quad C = O(1). \quad (4.13)$$

In view of assumption (1.4), we have $\mathcal{P}([\tilde{\alpha}, 1]) > 0$. Then, for any $r > 1 - \delta$ we have

$$\begin{aligned}
A &\geq \int_{\tilde{\alpha}}^1 \int_{\cup_{i=1}^m \varphi_{\theta_i}^{-1}(B_\delta(0))} e^{\alpha u_{r,\sigma}} \mathcal{P}(d\alpha) dx \\
&\geq \mathcal{P}([\tilde{\alpha}, 1]) \int_{\cup_{i=1}^m \varphi_{\theta_i}^{-1}(B_\delta(0))} e^{\tilde{\alpha} u_{r,\sigma}} dx \\
&= \mathcal{P}([\tilde{\alpha}, 1]) \int_{\cup_{i=1}^m \varphi_{\theta_i}^{-1}(B_\delta(0))} \sum_{j=1}^m t_j e^{\tilde{\alpha} v_{r,\theta_j}} dx \\
&\geq \mathcal{P}([\tilde{\alpha}, 1]) \sum_{j=1}^m t_j \int_{\varphi_{\theta_j}^{-1}(B_\delta(0))} e^{\tilde{\alpha} v_{r,\theta_j}} dx \\
&= \varepsilon_0^2 \mathcal{P}([\tilde{\alpha}, 1]) \sum_{j=1}^m t_j \left[\int_{B_{1-r}(0)} \frac{dy}{(1-r)^{4\tilde{\alpha}}} + \int_{B_\delta(0) \setminus B_{1-r}(0)} \frac{dy}{|y|^{4\tilde{\alpha}}} \right] \\
&= \pi \varepsilon_0^2 \mathcal{P}([\tilde{\alpha}, 1]) \left[(1-r)^{2-4\tilde{\alpha}} + \frac{1}{2\tilde{\alpha}-1} ((1-r)^{2-4\tilde{\alpha}} - \delta^{2-4\tilde{\alpha}}) \right] \rightarrow +\infty \quad \text{uniformly for } \sigma \in \Gamma_k.
\end{aligned}$$

In the last line we have used that $\sum_{j=1}^m t_j = 1$ and that $\tilde{\alpha} > \frac{1}{2}$.

Moreover

$$\begin{aligned}
0 \leq B \leq C &\leq |\Omega| + \int_{\cup_{i=1}^m (B_{\varepsilon_0}(\gamma(\theta_i)) \setminus \varphi_{\theta_i}^{-1}(B_\delta(0)))} \left(\sum_{j=1}^m t_j e^{\tilde{\alpha} v_{r,\theta_j}} \right)^{1/\tilde{\alpha}} dx \\
&\leq |\Omega| + \int_{\cup_{i=1}^m (B_{\varepsilon_0}(\gamma(\theta_i)) \setminus \varphi_{\theta_i}^{-1}(B_\delta(0)))} \left(\sum_{j=1}^m \frac{t_j}{\delta^{4\tilde{\alpha}}} \right)^{1/\tilde{\alpha}} dx \\
&\leq |\Omega| + k\pi \frac{\varepsilon_0^2}{\delta^4}.
\end{aligned}$$

Hence, (4.13) is established. Letting $r \rightarrow 1$ in (4.12), we obtain (4.11) and, in turn, (4.6). This concludes the proof. \square

Proof of Proposition 4.1. Let us consider a continuous function $\eta : [0, 1] \rightarrow [0, 1]$ such that $\eta([0, \frac{1}{3}]) = 0$ and $\eta([\frac{2}{3}, 1]) = 1$ and let us introduce the map $h : D_k \rightarrow H_0^1(\Omega)$ as

$$h(\underline{z}_k) = \eta(|\underline{z}_k|) u_{|\underline{z}_k|^2, \sigma(\underline{z}_k)}, \quad (4.14)$$

where $\sigma(\underline{z}_k) = \frac{\sum_{i=1}^k |z_i|^2 \delta_{\gamma(\theta_i)}}{|\underline{z}_k|^2}$ and $\underline{z}_k = (z_1, \dots, z_k) = (|z_1|e^{i\theta_1}, \dots, |z_k|e^{i\theta_k})$.

We claim that $h \in \mathcal{F}_\lambda$. Indeed, by means of (4.5)–(4.6) it is immediate to see that h satisfies property (4.2)–(i). Moreover,

$$d\mu(h(\underline{z}_k)) \rightarrow |z_1|^2 \delta_{\gamma(\theta_1)} + \dots + |z_k|^2 \delta_{\gamma(\theta_k)} \quad \text{as } \underline{z}_k \rightarrow \partial D_k,$$

which in turn this implies that

$$m \circ h(\underline{z}_k) \rightarrow \tilde{\Psi}_k(\underline{z}_k) \quad \text{as } \underline{z}_k \rightarrow \partial D_k,$$

where Φ is defined in (2.16), so that (4.2)–(ii) is also fulfilled. Finally, by Lemma 2.6, we also deduce property (4.2)–(iii). \square

We are now ready to define, for $\lambda \in (8k\pi, 8(k+1)\pi)$, the min-max value:

$$c_\lambda = \inf_{h \in \mathcal{F}_\lambda} \sup_{u \in h(D_k)} J_\lambda(u). \quad (4.15)$$

In view of Proposition 4.1, we have $c_\lambda < +\infty$. The following lower bound relies in an essential way on the non-contractibility of Ω as assumed in (1.9).

Proposition 4.3. *Assume (1.4)–(1.9). Let $k \in \mathbb{N}$ and $\lambda \in (8k\pi, 8(k+1)\pi)$, then $c_\lambda > -\infty$.*

Proof. The case $k = 1$ has been treated in [21], while the case $k > 1$ can be worked out following [2] with minor modifications. We prove it for reader's convenience.

We assume by contradiction that for any $n \in \mathbb{N}$ there exists $h_n \in \mathcal{F}_\lambda$ such that

$$\sup_{u \in h_n(D_k)} J_\lambda(u) \leq -n.$$

In view of property (4.2)–(iii) in the definition of \mathcal{F} , for any $n \in \mathbb{N}$ we can find $u_n \in h_n(D_k) \subset H_0^1(\Omega)$ such that

$$J_\lambda(u_n) \leq -n, \quad \text{and} \quad m(u_n) = \underline{0}_k.$$

Next we can apply Lemma 2.5 with r and ε to be chosen later in a convenient way. Denoting by ℓ the positive integer (less or equal than k) found in the above mentioned Lemma, for any $j \in \{1, \dots, \ell\}$

$$\begin{aligned} 0 = m_j(u_n) &= \int_{\Omega} (\chi(x))^j d\mu(u_n) = \int_{\cup_{i=1}^{\ell} B_r(p_i)} (\chi(x))^j d\mu(u_n) + \int_{\Omega \setminus \cup_{i=1}^{\ell} B_r(p_i)} (\chi(x))^j d\mu(u_n) \\ &= \sum_{i=1}^{\ell} \int_{B_r(p_i) \setminus \cup_{h=1}^{i-1} B_r(p_h)} (\chi(x))^j d\mu(u_n) + \int_{\Omega \setminus \cup_{i=1}^{\ell} B_r(p_i)} (\chi(x))^j d\mu(u_n) \\ &= \sum_{i=1}^{\ell} \beta_i (\chi(p_i))^j - R_{j,n}(r) \end{aligned} \quad (4.16)$$

where β_i and p_i are obtained via Lemma 2.5 and $R_{j,n}(r)$, up to a subsequence, can be estimated as follows:

$$\begin{aligned} |R_{j,n}(r)| &= \left| \sum_{i=1}^{\ell} \left(\int_{B_r(p_i) \setminus \cup_{h=1}^{i-1} B_r(p_h)} (\chi(x))^j d\mu(u_n) - (\chi(p_i))^j \beta_i \right) + \int_{\Omega \setminus \cup_{i=1}^{\ell} B_r(p_i)} (\chi(x))^j d\mu(u_n) \right| \\ &\stackrel{(2.15)+(2.14)}{\leq} \sum_{i=1}^{\ell} \int_{B_r(p_i) \setminus \cup_{h=1}^{i-1} B_r(p_h)} |(\chi(x))^j - (\chi(p_i))^j| d\mu(u_n) + o_n(1) + \varepsilon \\ &\leq \sum_{i=1}^{\ell} \int_{B_r(p_i) \setminus \cup_{h=1}^{i-1} B_r(p_h)} |\chi(x) - \chi(p_i)| \left(\sum_{h=0}^{j-1} |\chi(x)|^{j-1-h} |\chi(p_i)|^h \right) d\mu(u_n) + o_n(1) + \varepsilon \\ &\leq \sum_{i=1}^{\ell} j d^{j-1} \int_{B_r(p_i) \setminus \cup_{h=1}^{i-1} B_r(p_h)} |\chi(x) - \chi(p_i)| d\mu(u_n) + o_n(1) + \varepsilon \\ &\leq \sum_{i=1}^{\ell} \ell d^{j-1} C_\chi r \beta_i + o_n(1) + \varepsilon = \ell d^{j-1} C_\chi r + o_n(1) + \varepsilon. \end{aligned}$$

In the above chain of inequalities $d := \max_{x \in \Omega} |\chi(x)|$ and $C_\chi = \max_{x_1, x_2 \in \Omega} \frac{|\chi(x_1) - \chi(x_2)|}{|x_1 - x_2|}$. Denoting by $z_i := \chi(p_i) \in \mathbb{C}$, for $i \in \{1, \dots, \ell\}$, we get, by virtue of (4.16), that the z_i 's satisfy

$$\begin{cases} \beta_1 z_1 + \beta_2 z_2 + \dots + \beta_\ell z_\ell = R_{1,n}(r) \\ \beta_1 z_1^2 + \beta_2 z_2^2 + \dots + \beta_\ell z_\ell^2 = R_{2,n}(r) \\ \dots\dots\dots \\ \beta_1 z_1^\ell + \beta_2 z_2^\ell + \dots + \beta_\ell z_\ell^\ell = R_{\ell,n}(r). \end{cases} \quad (4.17)$$

By our choice of χ , see (4.1), there exists $\rho > 0$ such that $\chi(\Omega) \cap B_{2\rho}(0) = \emptyset$, then

$$2\rho \leq |\chi(p_i)| = |z_i|. \quad (4.18)$$

On the other hand, by applying Lemma 2.7 to system (4.17), we obtain that there exists $\delta > 0$ such that if

$$|R_{j,n}(r)| \leq \delta \quad \text{for some } n \in \mathbb{N} \text{ and for any } j \in \{1, \dots, \ell\}, \quad (4.19)$$

then $|z_i| \leq \rho$, which would be a contradiction against (4.18). Finally, it is immediate to see that choosing $r = \frac{\delta}{2\ell d^{j-1} C_\chi}$, $\varepsilon = \frac{\delta}{2}$ and n sufficiently large condition (4.19) is fulfilled for any $j \in \{1, \dots, \ell\}$. The proof is thereby complete. \square

Finally, we are able to prove the existence result.

Proof of Theorem 1.2. By the definition (1.5) of J_λ , it is readily checked that if $\lambda' \leq \lambda$, then $\mathcal{F}_{\lambda'} \subset \mathcal{F}_\lambda$ and $J_{\lambda'}(u) \geq J_\lambda(u)$ for all $u \in H_0^1(\Omega)$. Consequently, $c_{\lambda'} \geq c_\lambda$, where c_λ is the min-max value defined in (4.15). In particular, the mapping $\lambda \rightarrow c_\lambda$ is monotone, and therefore the derivative c'_λ exists for almost every $\lambda \in (8k\pi, 8(k+1)\pi)$. We fix $\lambda \in (8k\pi, 8(k+1)\pi)$ such that c'_λ is well-defined. By the well-known Struwe Monotonicity Trick [24], a bounded Palais-Smale sequence, whose bounds depend on $|c'_\lambda|$, may be constructed at level c_λ . The details of this construction in the context of mean field equations may be found in [25], see also [21] for the specific context of $(*)_\lambda$. By compactness of the Moser-Trudinger embedding, we obtain from the bounded Palais-Smale sequence a solution to $(*)_\lambda$. In this way, we obtain a solution to $(*)_\lambda$ for almost every $\lambda \in (8k\pi, 8(k+1)\pi)$. Now we fix $\lambda_0 \in (8k\pi, 8(k+1)\pi)$. Let $\lambda_n \rightarrow \lambda_0$ be such that $(*)_{\lambda_n}$ admits a solution u_n for all n . In view of Theorem 1.1, we conclude that the sequence u_n is compact, and consequently there exists a solution u_0 to $(*)_{\lambda_0}$ such that $u_n \rightarrow u_0$. In particular, we obtain a solution for $(*)_{\lambda_0}$. We conclude that solutions to $(*)_\lambda$ exist for all values $\lambda \in (8k\pi, 8(k+1)\pi)$, as asserted. \square

5. APPENDIX: LIOUVILLE BUBBLE LIMIT PROFILES

In view of the mass quantization property, as stated in Theorem 1.1, it is natural to expect that, upon rescaling, a concentrating sequence u_n of solutions to $(*)_{\lambda_n}$ should yield a *Liouville bubble* profile, namely a profile of the form

$$U_{\delta, \xi}(x) = \log \frac{8\delta^2}{(\delta^2 + |x - \xi|^2)^2}, \quad \delta > 0, \xi \in \mathbb{R}^2. \quad (5.1)$$

This is indeed the case, as we show in this Appendix. However, it turns out that the usual rescaling yields the desired profile only in the “non-degenerate” case where $\mathcal{P}(\{1\}) > 0$. On the other hand, if $\mathcal{P}(\{1\}) = 0$, such a rescaling yields a trivial profile in the limit, and some extra care is needed in order to capture the Liouville bubble profile.

More precisely, let u_n be a concentrating sequence of solutions to $(*)_{\lambda_n}$. It is convenient to set

$$\mathcal{I}_n = \iint_{[0,1] \times \Omega} e^{\alpha' u_n} \mathcal{P}(d\alpha') dx.$$

We recall that by the maximum principle $\mathcal{I}_n \geq |\Omega|$, and along a concentrating sequence we have $\mathcal{I}_n \rightarrow +\infty$. Then, problem $(*)_{\lambda_n}$ takes the form

$$\begin{cases} -\Delta u_n = \frac{\lambda_n}{\mathcal{I}_n} \int_{[0,1]} \alpha e^{\alpha u_n} \mathcal{P}(d\alpha) & \text{in } \Omega \\ u_n = 0 & \text{on } \partial\Omega. \end{cases}$$

We assume that $\lambda_n \rightarrow \lambda_0$ and that

$$u_n(x_n) = \max_{\Omega} u_n \rightarrow +\infty \quad \text{as } n \rightarrow \infty.$$

In view of [12] (see the proof of Lemma 2.1), we know that x_n stays well-away from $\partial\Omega$.

5.1. The “non-degenerate case” $\mathcal{P}(\{1\}) > 0$. Throughout this subsection we assume that

$$1 \in \text{supp } \mathcal{P} \quad \text{and} \quad \mathcal{P}(\{1\}) > 0.$$

We define

$$w_n(x) := u_n(x) - \log \mathcal{I}_n.$$

Then, $w_n(x_n) = \max w_n$ and w_n satisfies

$$-\Delta w_n = \frac{\lambda_n}{\mathcal{I}_n} \int_{[0,1]} \alpha e^{\alpha u_n} \mathcal{P}(d\alpha) = \lambda_n \mathcal{P}(\{1\}) e^{w_n} + \rho_n,$$

where

$$\rho_n(x) := \frac{\lambda_n}{\mathcal{I}_n} \int_{[0,1]} \alpha e^{\alpha u_n} \mathcal{P}(d\alpha).$$

Moreover,

$$\begin{aligned} \int_{\Omega} e^{w_n} dx &= \int_{\Omega} e^{u_n - \log \mathcal{I}_n} dx = \int_{\Omega} \frac{e^{u_n} dx}{\iint_{[0,1] \times \Omega} e^{\alpha' u_n} \mathcal{P}(d\alpha') dx} \\ &= \mathcal{P}(\{1\})^{-1} \frac{\iint_{\{1\} \times \Omega} e^{\alpha' u_n} \mathcal{P}(d\alpha') dx}{\iint_{[0,1] \times \Omega} e^{\alpha' u_n} \mathcal{P}(d\alpha') dx} \\ &\leq \mathcal{P}(\{1\})^{-1}. \end{aligned}$$

In order to rescale, we set:

$$\sigma_n := e^{-w_n(x_n)/2} \quad \tilde{\Omega}_n := \frac{\Omega - x_n}{\sigma_n}$$

and

$$\tilde{w}_n(y) := w_n(x_n + \sigma_n y) + 2 \log \sigma_n, \quad y \in \tilde{\Omega}_n.$$

Since x_n stays well-away from the boundary of Ω , the rescaled domain $\tilde{\Omega}_n$ invades the whole space \mathbb{R}^2 .

Then, \tilde{w}_n satisfies

$$\begin{cases} -\Delta \tilde{w}_n = \lambda_n \mathcal{P}(\{1\}) e^{\tilde{w}_n} + \lambda_n \tilde{\rho}_n & \text{in } \tilde{\Omega}_n \\ \tilde{w}_n(y) \leq \tilde{w}_n(0) = 0 \\ \int_{\tilde{\Omega}_n} e^{\tilde{w}_n} = \int_{\Omega} e^{w_n} \leq \mathcal{P}(\{1\})^{-1} \end{cases}$$

where

$$\tilde{\rho}_n(y) := \frac{\sigma_n^2}{\mathcal{I}_n} \int_{[0,1]} \alpha e^{\alpha u_n(x_n + \sigma_n y)} \mathcal{P}(d\alpha).$$

We claim that

$$\|\tilde{\rho}_n\|_{L^\infty(\tilde{\Omega}_n)} \rightarrow 0 \text{ as } n \rightarrow \infty. \quad (5.2)$$

Indeed, given $\varepsilon > 0$, let $\eta > 0$ be sufficiently small so that $\mathcal{P}([1-\eta, 1]) < \varepsilon$. Let $n_0 \in \mathbb{N}$ be sufficiently large so that $(\sigma_n^2/\mathcal{I}_n)^\eta < \varepsilon$ for all $n \geq n_0$. We estimate, for all $n \geq n_0$:

$$\begin{aligned} \tilde{\rho}_n(y) &= \int_{[0,1]} \alpha e^{\alpha \tilde{w}_n(y)} \left(\frac{\sigma_n^2}{\mathcal{I}_n} \right)^{1-\alpha} \mathcal{P}(d\alpha) \stackrel{\tilde{w}_n(y) \leq 0}{\leq} \int_{[0,1]} \left(\frac{\sigma_n^2}{\mathcal{I}_n} \right)^{1-\alpha} \mathcal{P}(d\alpha) \\ &= \int_{[0,1-\eta]} \left(\frac{\sigma_n^2}{\mathcal{I}_n} \right)^{1-\alpha} \mathcal{P}(d\alpha) + \int_{[1-\eta,1]} \left(\frac{\sigma_n^2}{\mathcal{I}_n} \right)^{1-\alpha} \mathcal{P}(d\alpha) \\ &\leq \left(\frac{\sigma_n^2}{\mathcal{I}_n} \right)^\eta + \mathcal{P}([1-\eta, 1]) < 2\varepsilon. \end{aligned}$$

Hence, (5.2) is established.

We conclude that there exists a solution $\tilde{w} \in C_{\text{loc}}^2(\mathbb{R}^2)$ to the problem

$$\begin{cases} -\Delta \tilde{w} = \lambda_0 \mathcal{P}(\{1\}) e^{\tilde{w}} & \text{in } \mathbb{R}^2 \\ \tilde{w}(y) \leq \tilde{w}(0) = 0 \\ \int_{\mathbb{R}^2} e^{\tilde{w}} \leq \mathcal{P}(\{1\})^{-1} \end{cases}$$

such that, up to subsequences, $\tilde{w}_n \rightarrow \tilde{w}$ in $C_{\text{loc}}^2(\mathbb{R}^2)$. In view of Chen-Li's classification result [6], the function $\tilde{w} + \log(\lambda_0 \mathcal{P}(\{1\}))$ is of the form (5.1). The asserted limit profile is thus established in the case $\mathcal{P}(\{1\}) > 0$.

5.2. The “degenerate case” $\mathcal{P}(\{1\}) = 0$. Throughout this section, we assume

$$1 \in \text{supp } \mathcal{P} \quad \text{and} \quad \mathcal{P}(\{1\}) = 0. \quad (5.3)$$

We show the following.

Proposition 5.1. *Assume (5.3). There exist a rescaling of u_n of the form*

$$\tilde{w}_n(y) = \alpha_n u_n(x_n + \sigma_n y) - \log \mathcal{I}_n, \quad y \in \tilde{\Omega}_n = \frac{\Omega - x_n}{\sigma_n}$$

where, as $n \rightarrow \infty$, $\alpha_n \rightarrow 1$, $\sigma_n^2 = e^{-\alpha_n u_n(x_n) + \log \mathcal{I}_n} \rightarrow 0$, and a solution \tilde{w} to the problem

$$\begin{cases} -\Delta \tilde{w} = \lambda_0 e^{\tilde{w}} & \text{in } \mathbb{R}^2 \\ \int_{\mathbb{R}^2} e^{\tilde{w}} dx < +\infty \\ \tilde{w}(y) \leq \tilde{w}(0) = 0, \end{cases}$$

such that, up to subsequences, $\tilde{w}_n \rightarrow \tilde{w}$ in $C_{\text{loc}}^2(\mathbb{R}^2)$. In particular, $\tilde{w} + \log \lambda_0$ is of the desired form (5.1).

We define $\alpha_n \in [0, 1]$ and the functions w_n, V_n by setting:

$$\begin{aligned} e^{\alpha_n u_n(x_n)} &:= \int_{[0,1]} \alpha e^{\alpha u_n(x_n)} \mathcal{P}(d\alpha) \\ w_n &:= \alpha_n u_n - \log \mathcal{I}_n \\ V_n &:= \frac{\alpha_n \lambda_n}{\mathcal{I}_n} \int_{[0,1]} \alpha e^{\alpha u_n} \mathcal{P}(d\alpha) e^{-w_n}. \end{aligned}$$

With this notation, we have:

Lemma 5.2. *Assume (5.3). The following facts hold true:*

- (i) $\alpha_n \rightarrow 1$.
- (ii) $V_n(x_n) = \alpha_n \lambda_n \rightarrow \lambda_0$
- (iii) $\|V_n\|_{L^\infty(\Omega)} \leq \alpha_n \lambda_n (\int_{[0,1]} \alpha \mathcal{P}(d\alpha) + 1)$
- (iv) $\int_{\Omega} V_n e^{w_n} dx \leq \alpha_n \lambda_n$.

Proof. Proof of (i). There holds:

$$e^{\alpha_n} = \left(\int_{[0,1]} \alpha e^{\alpha u_n(x_n)} \mathcal{P}(d\alpha) \right)^{1/u_n(x_n)} = \|\alpha e^\alpha\|_{L^{u_n(x_n)}(I, \mathcal{P})} \rightarrow \|\alpha e^\alpha\|_{L^\infty(I, \mathcal{P})} = e.$$

Hence, $\alpha_n \rightarrow 1$.

Proof of (ii). By definition of α_n and w_n , we have

$$V_n(x_n) = \frac{\alpha_n \lambda_n}{\mathcal{I}_n} \int_{[0,1]} \alpha e^{\alpha u_n(x_n)} \mathcal{P}(d\alpha) e^{-w_n(x_n)} = \frac{\alpha_n \lambda_n}{\mathcal{I}_n} e^{\alpha_n u_n(x_n)} e^{-w_n(x_n)} = \alpha_n \lambda_n.$$

Proof of (iii). By definition, we have

$$V_n = \alpha_n \lambda_n \int_{[0,1]} \alpha e^{(\alpha - \alpha_n) u_n} \mathcal{P}(d\alpha).$$

Since $u_n \geq 0$, we estimate:

$$\begin{aligned} \int_{[0,1]} \alpha e^{(\alpha - \alpha_n) u_n} \mathcal{P}(d\alpha) &= \int_{\alpha < \alpha_n} \alpha e^{(\alpha - \alpha_n) u_n} \mathcal{P}(d\alpha) + \int_{\alpha \geq \alpha_n} \alpha e^{(\alpha - \alpha_n) u_n} \mathcal{P}(d\alpha) \\ &\leq \int_{\alpha < \alpha_n} \alpha \mathcal{P}(d\alpha) + \int_{\alpha \geq \alpha_n} \alpha e^{(\alpha - \alpha_n) u_n(x_n)} \mathcal{P}(d\alpha) \\ &\leq \int_{[0,1]} \alpha \mathcal{P}(d\alpha) + \int_{[0,1]} \alpha e^{(\alpha - \alpha_n) u_n(x_n)} \mathcal{P}(d\alpha) \\ &= \int_{[0,1]} \alpha \mathcal{P}(d\alpha) + 1. \end{aligned}$$

The asserted estimate follows.

Proof of (iv). By definition of \mathcal{I}_n, V_n , we have

$$\int_{\Omega} V_n e^{w_n} = \frac{\alpha_n \lambda_n}{\mathcal{I}_n} \iint_{I \times \Omega} \alpha e^{\alpha u_n} \mathcal{P}(d\alpha) \leq \alpha_n \lambda_n.$$

The asserted estimates are established. \square

Now, we can prove Proposition 5.1.

Proof of Proposition 5.1. We define the rescaling:

$$\sigma_n := e^{w_n(x_n)/2}, \quad \tilde{w}_n(y) := w_n(x_n + \sigma_n y) + 2 \log \sigma_n, \quad \tilde{V}_n(y) = V_n(x_n + \sigma_n y).$$

The function \tilde{w}_n satisfies

$$\begin{cases} -\Delta \tilde{w}_n = \tilde{V}_n e^{\tilde{w}_n} & \text{in } \tilde{\Omega}_n \\ \int_{\tilde{\Omega}_n} \tilde{V}_n e^{\tilde{w}_n} \leq C \\ \tilde{w}_n(y) \leq \tilde{w}_n(0) = 0 \\ \|\tilde{V}_n\|_{L^\infty(\tilde{\Omega}_n)} \leq C, \end{cases}$$

where, as above, $\tilde{\Omega}_n$ invades the whole space \mathbb{R}^2 . In view of the estimates in Lemma 5.2, there exists a solution \tilde{w} to the problem

$$\begin{cases} -\Delta \tilde{w} = \lambda_0 e^{\tilde{w}} \\ \int_{\mathbb{R}^2} e^{\tilde{w}} < +\infty \\ \tilde{w}(y) \leq \tilde{w}(0) = 0. \end{cases}$$

such that a subsequence, still denoted \tilde{w}_n , satisfies $\tilde{w}_n \rightarrow \tilde{w}$ locally uniformly on \mathbb{R}^2 . In view of the classification in [6], the function $\tilde{w} + \log \lambda_0$ is of the form (5.1).

Hence, Proposition 5.1 is established. \square

In view of Proposition 5.1, we expect that concentrating solutions to $(*)_\lambda$ may be constructed by the approach in [11].

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REFERENCES

- [1] Aubin, T., *Some Nonlinear Problems in Riemannian Geometry*, Springer Monographs in Mathematics, Springer-Verlag, Berlin, 1998.
- [2] Bartolucci, D., De Marchis, F., *On the Ambjorn-Olesen electroweak condensates*, J. Math. Phys. **53** (2012), no. 7, 073704, 15 pp.
- [3] Bartolucci, D., De Marchis, F., *Supercritical mean field equations on convex domains and the Onsager's statistical description of two-dimensional turbulence*, Arch. Ration. Mech. Anal. **217** (2015), no. 2, 525–570.
- [4] Brezis, H., Merle, F., *Uniform estimates and blow-up behavior for solutions of $-\Delta u = V(x)e^u$ in dimension 2*, Comm. Partial Differential Equations **16** (1991), 1223–1253.
- [5] Caglioti, E., Lions, P.L., Marchioro, C., Pulvirenti, M., *A special class of stationary flows for two-dimensional Euler equations: a statistical mechanics description*, Comm. Math. Phys. **143** (1992), no. 3, 501–525.
- [6] Chen, W., Li, C., *Classification of solutions of some nonlinear elliptic equations*, Duke Math. J. **63** no. 3 (1991), 615–622.
- [7] Chen, W., Li, C., *Prescribing Gaussian Curvatures on Surfaces with Conical Singularities*, J. of Geometric Analysis **1** n. 4 (1991), 359–372.
- [8] Ding, W., Jost, J., Li, J., Wang, G., *Existence results for mean field equations*, Ann. Inst. H. Poincaré, **16** no. 5 (1999), 653–666.
- [9] Djadli, Z., *Existence result for the mean field problem on Riemann surfaces of all genres*, Commun. Contemp. Math. **10** (2008), 205–220.
- [10] Djadli, Z., Malchiodi, A., *Existence of conformal metrics with constant Q -curvature*, Ann. of Math. **168** no. 3 (2008), 813–858.
- [11] Esposito, P., Grossi, M., Pistoia, A., *On the existence of blowing-up solutions for a mean field equation*, Ann. I. H. Poincaré - AN **22** (2005), 227–257.
- [12] Gidas, B. Ni, W.M., Nirenberg, L., *Symmetry and Related Properties via the Maximum Principle*, Commun. Math. Phys. **68** (1979), 209–243.
- [13] Lin, C.S., *An expository survey on recent development of mean field equations*, Discr. Cont. Dynamical Systems **19** n. 2 (2007), 217–247.
- [14] Malchiodi, A., *Topological methods for an elliptic equation with exponential nonlinearities*, Discrete Cont. Dyn. Syst. **21** (2008), 277–294.
- [15] Moser, J. *A sharp form of an inequality by N. Trudinger*, Indiana Univ. Math. J. **20** (1970/71), 1077–1092.

- [16] Nagasaki, K., Suzuki, T., *Asymptotic analysis for two-dimensional elliptic eigenvalue problems with exponentially-dominated nonlinearities*, Asymptotic Analysis 3 (1990), 173–188.
- [17] Neri, C. *Statistical Mechanics of the N-point vortex system with random intensities on a bounded domain*, Ann. I. H. Poincaré - AN **21** (2004), 381–399. J. Differential Equations 249 n. 6 (2010), 1436–1465.
- [18] Onsager, L., *Statistical hydrodynamics*, Nuovo Cimento Suppl. n. 2 **6** (9), (1949), 279–287.
- [19] Ricciardi, T., Zecca, G., *Blow-up analysis for some mean field equations involving probability measures from statistical hydrodynamics*, Differential and Integral Equations 25 n. 3–4 (2012), 201–222.
- [20] Ricciardi, T., Zecca, G., *On the blow-up of solutions to Liouville type equations*, Adv. Nonlinear Stud. **16** no. 1 (2016), 75–85.
- [21] Ricciardi, T., Zecca, G., *Mass quantization and minimax solutions for Neri’s mean field equation in 2D-turbulence*, J. Differential Equations **260** (2016), 339–369.
- [22] Schoenflies, A., *Beiträge zur Theorie der Punktmengen. III.*, Math. Ann. **62** (1906), 286–328.
- [23] Senba, T., Suzuki, T., *Chemotactic collapse in a parabolic-elliptic system of mathematical biology*, Adv. Differential Equations **6** (2001), 21–50.
- [24] Struwe, M., *The existence of surfaces of constant mean curvature with free boundaries*, Acta Math. **160** (1988), 19–64.
- [25] Struwe, M., Tarantello, G., *On multivortex solutions in Chern-Simons gauge theory*, Boll. Unione Mat. Ital. Sez.B 1 (1998), 109–121.
- [26] Suzuki, T., *Mean Field Theories and Dual Variation*, Atlantis Press, Paris, 2009.
- [27] Trudinger, N.S., *On imbeddings into Orlicz spaces and some applications*, J. Math. Mech. **17** (1967), 473–483.
- [28] Ye, D., *Une remarque sur le comportement asymptotique des solutions de $-\Delta u = \lambda f(u)$* , C.R. Acad. Sci. Paris **325** (1997), 1279–1282.

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